

Last time we tried to build an invariant  $\tilde{\psi}\psi$  from spinors trying  $\tilde{\psi} = \psi^*$  (like we did with  $SU(2)$ ),

but...  $\tilde{\psi}^T \psi = (\psi^*)^T \psi = \psi^\dagger \psi \rightarrow (\psi')^\dagger \psi' = (e^{\frac{i}{4} \sigma^{\mu\nu} \omega_{\mu\nu}} \psi)^\dagger e^{\frac{i}{4} \sigma^{\mu\nu} \omega_{\mu\nu}} \psi$

Recall:  $\sigma^{0i} = \frac{i}{2} \begin{pmatrix} \sigma_i & 0 \\ 0 & -\sigma_i \end{pmatrix}$   $\sigma^{ij} = \frac{i}{2} \epsilon^{ijk} \begin{pmatrix} \sigma_k & 0 \\ 0 & \sigma_k \end{pmatrix}$   
and  $\sigma^{i\dagger} = \sigma^i$

$= \psi^\dagger (e^{\frac{i}{4} \sigma^{\mu\nu} \omega_{\mu\nu}})^\dagger e^{\frac{i}{4} \sigma^{\mu\nu} \omega_{\mu\nu}} \psi$

But the  $\sigma^{\mu\nu}$  are not all Hermitian!

In particular  $\sigma^{0i\dagger} = -\sigma^{0i}$  Not-Hermitian  
 $\sigma^{ij\dagger} = \sigma^{ij}$  Hermitian

The reason  $\psi^\dagger \psi$  worked for  $SU(2)$  is that the generators were all Hermitian.

But we can fix this with a different choice of dual:  $\tilde{\psi} = i\gamma^0 \psi^*$

Then:  $\tilde{\psi}^T \psi = (i\gamma^0 \psi^*)^T \psi$

$= i\psi^\dagger \gamma^0 \psi$

$= i\psi^\dagger \gamma^0 \psi \rightarrow i(\psi')^\dagger \gamma^0 \psi' = i(e^{\frac{i}{4} \sigma^{\mu\nu} \omega_{\mu\nu}} \psi)^\dagger \gamma^0 e^{\frac{i}{4} \sigma^{\mu\nu} \omega_{\mu\nu}} \psi$

$= i\psi^\dagger (e^{\frac{i}{4} \sigma^{\mu\nu} \omega_{\mu\nu}})^\dagger \gamma^0 e^{\frac{i}{4} \sigma^{\mu\nu} \omega_{\mu\nu}} \psi$   
 $\gamma^0 e^{-\frac{i}{4} \sigma^{\mu\nu} \omega_{\mu\nu}}$

You will show this in your HW

Same!  $\rightarrow = i\psi^\dagger \gamma^0 \psi$

So in the end we define our dual spinor with  $\tilde{\psi} = i\gamma^0 \psi^*$  and the adjoint  $\bar{\psi} \equiv \tilde{\psi}^T = i\psi^\dagger \gamma^0$

Or in other words  $g = i\gamma^0 \Rightarrow \tilde{\psi} = g\psi$

$\Downarrow$   
 $\bar{\psi}\psi$  is invariant